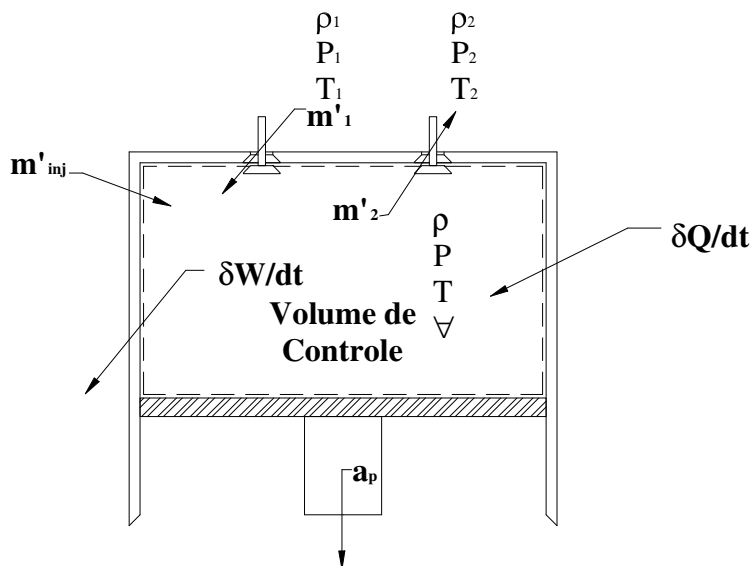


ANEXO D

DEDUÇÕES ALGÉBRICAS:

D.1 MODELAGEM MATEMÁTICA ZERO-DIMENSIONAL DE UMA ZONA, PARA UM CILINDRO DE MOTOR A PISTÃO



Equação da Conservação da Massa:

$$\iint_{S.C.} \rho \vec{V} \cdot d\vec{A} = -\frac{\partial}{\partial t} \iiint_{V.C.} \rho dV$$

$$\iint_{A_1} \rho_1 \vec{V}_1 \cdot d\vec{A} + \iint_{A_2} \rho_2 \vec{V}_2 \cdot d\vec{A} = -\frac{\partial}{\partial t} \iiint_{V.C.} \rho dV$$

Integrando ao longo do volume de controle e considerando uma distribuição homogênea no interior do mesmo, tem-se,

$$-\dot{m}_{inj} - \dot{m}_1 + \dot{m}_2 = -\frac{d}{dt}(\rho V)$$

onde, o sub-índice 1 é para o fluxo que entra no volume de controle e o sub-índice 2 é para o fluxo que deixa o volume de controle. A vazão mássica de injeção direta de combustível no cilindro é representada por

$$\dot{m}_{inj}$$

Mas a massa específica do gás que sai é igual a do gás no interior do cilindro, supondo uma distribuição homogênea dentro do mesmo, assim,

$$\rho_2 = \rho$$

$$\rho \frac{d\mathcal{V}}{dt} + \mathcal{V} \frac{d\rho}{dt} = \dot{m}_1 + \dot{m}_{inj} - \dot{m}_2$$

$$\mathcal{V} \frac{d\rho}{dt} = \dot{m}_1 + \dot{m}_{inj} - \dot{m}_2 - \rho \frac{d\mathcal{V}}{dt}$$

$$\frac{d\rho}{dt} = \frac{1}{\mathcal{V}} \left(\dot{m}_1 + \dot{m}_{inj} - \dot{m}_2 - \rho \frac{d\mathcal{V}}{dt} \right)$$

Mudando da variável independente tempo para posição angular do virabrequim (θ), tem-se,

$$\frac{d\theta}{dt} = \omega \Rightarrow dt = \frac{d\theta}{\omega}$$

onde, ω = velocidade de rotação (rad/s)

Portanto,

$$\frac{d\rho}{d\theta} = \frac{1}{\omega \mathcal{V}} \left(\dot{m}_1 + \dot{m}_{inj} - \dot{m}_2 - \rho \omega \frac{d\mathcal{V}}{d\theta} \right)$$

Equação da Conservação de Energia:

$$\frac{\delta Q}{dt} - \frac{\delta W_s}{dt} = \iint_{s.c.} \left(e + \frac{P}{\rho} \right) \rho (\vec{V} \cdot \vec{n}) dA + \frac{\partial}{\partial t} \iiint_{v.c.} \rho e p d\mathcal{V} + \frac{\delta W_\mu}{dt}$$

onde, $e = u + \frac{V^2}{2} + gz$

para o motor alternativo,

$$\frac{\delta W_s}{dt} = P \frac{d\mathcal{V}}{dt}$$

$$\frac{\delta Q}{dt} = \frac{\delta Q_f}{dt} + \frac{\delta Q_1}{dt}$$

portanto,

$$\frac{\delta Q_f}{dt} + \frac{\delta Q_1}{dt} - P \frac{dV}{dt} = - \left(e_1 + \frac{P_1}{\rho_1} \right) \dot{m}_1 + \left(e_2 + \frac{P_2}{\rho_2} \right) \dot{m}_2 + \frac{d}{dt} (e \rho V) + \frac{\delta W_\mu}{dt}$$

$$\text{mas, } e + \frac{P}{\rho} = h + \frac{V^2}{2} + gz$$

$$\begin{aligned} V \frac{d(\rho e)}{dt} + \rho e \frac{dV}{dt} &= \frac{\delta Q_f}{dt} + \frac{\delta Q_1}{dt} - P \frac{dV}{dt} + \frac{d}{dt} (e \rho V) - \frac{\delta W_\mu}{dt} + \\ &+ \dot{m}_1 \left(h_1 + \frac{V_1^2}{2} + gz_1 \right) - \dot{m}_2 \left(h_2 + \frac{V_2^2}{2} + gz_2 \right) \end{aligned}$$

$$\begin{aligned} \rho V \frac{de}{dt} + e V \frac{d\rho}{dt} + \rho e \frac{dV}{dt} &= \frac{\delta Q_f}{dt} + \frac{\delta Q_1}{dt} - P \frac{dV}{dt} - \frac{\delta W_\mu}{dt} + \\ &+ \dot{m}_1 \left(h_1 + \frac{V_1^2}{2} + gz_1 \right) - \dot{m}_2 \left(h_2 + \frac{V_2^2}{2} + gz_2 \right) \end{aligned}$$

$$\begin{aligned} \rho V \frac{d}{dt} \left(u + \frac{V^2}{2} + gz \right) &= - \left(u + \frac{V^2}{2} + gz \right) \left(V \frac{d\rho}{dt} + \rho \frac{dV}{dt} \right) + \frac{\delta Q_f}{dt} + \frac{\delta Q_1}{dt} - P \frac{dV}{dt} - \frac{\delta W_\mu}{dt} + \\ &+ \dot{m}_1 \left(h_1 + \frac{V_1^2}{2} + gz_1 \right) - \dot{m}_2 \left(h_2 + \frac{V_2^2}{2} + gz_2 \right) \end{aligned}$$

$$\begin{aligned} \rho V \frac{du}{dt} + \rho V V \frac{dV}{dt} + \rho V \frac{d(gz)}{dt} &= - \left(h + \frac{V^2}{2} + gz - \frac{P}{\rho} \right) \left(V \frac{d\rho}{dt} + \rho \frac{dV}{dt} \right) + \frac{\delta Q_f}{dt} + \frac{\delta Q_1}{dt} - P \frac{dV}{dt} - \\ &- \frac{\delta W_\mu}{dt} + \dot{m}_1 \left(h_1 + \frac{V_1^2}{2} + gz_1 \right) - \dot{m}_2 \left(h_2 + \frac{V_2^2}{2} + gz_2 \right) \end{aligned}$$

$$\begin{aligned} \rho V c_v \frac{dT}{dt} &= - \rho V \left(V \frac{dV}{dt} + \frac{d(gz)}{dt} \right) - \left(h + \frac{V^2}{2} + gz - \frac{P}{\rho} \right) \left(V \frac{d\rho}{dt} + \rho \frac{dV}{dt} \right) + \frac{\delta Q_f}{dt} + \frac{\delta Q_1}{dt} - P \frac{dV}{dt} - \\ &- \frac{\delta W_\mu}{dt} + \dot{m}_1 \left(h_1 + \frac{V_1^2}{2} + gz_1 \right) - \dot{m}_2 \left(h_2 + \frac{V_2^2}{2} + gz_2 \right) \end{aligned}$$

$$\text{mas, } c_v = \frac{R}{(\gamma-1)} \quad (\text{para gás ideal}) \text{ e desprezando a energia potencial,}$$

$$\frac{dT}{dt} = \frac{(\gamma-1)}{\rho \Psi R} \left[-\rho \Psi \left(V \frac{dV}{dt} \right) - \left(h + \frac{V^2}{2} - \frac{P}{\rho} \right) \left(\Psi \frac{d\rho}{dt} + \rho \frac{d\Psi}{dt} \right) + \frac{\delta Q_f}{dt} + \frac{\delta Q_l}{dt} - P \frac{d\Psi}{dt} - \frac{\delta W_\mu}{dt} + \dot{m}_1 \left(h_1 + \frac{V_1^2}{2} \right) - \dot{m}_2 \left(h + \frac{V_2^2}{2} \right) \right]$$

trocando a variável independente,

$$\frac{dT}{d\theta} = \frac{(\gamma-1)}{\rho \Psi \omega R} \left[-\rho \Psi \omega \left(V \frac{dV}{d\theta} \right) - \omega \left(h + \frac{V^2}{2} - \frac{P}{\rho} \right) \left(\Psi \frac{d\rho}{d\theta} + \rho \frac{d\Psi}{d\theta} \right) + \omega \frac{\delta Q_f}{d\theta} + \omega \frac{\delta Q_l}{d\theta} - \omega P \frac{d\Psi}{d\theta} - \omega \frac{\delta W_\mu}{d\theta} + \dot{m}_1 \left(h_1 + \frac{V_1^2}{2} \right) + \dot{m}_{inj} h_f - \dot{m}_2 \left(h + \frac{V_2^2}{2} \right) \right]$$

Equação da Conservação da Quantidade de Movimento:

$$\Sigma \vec{F}_{ext} = \iint_{S.C.} \vec{V} \rho (\vec{V} \cdot \vec{n}) dA + \frac{\partial}{\partial t} \iiint_{V.C.} \rho \vec{V} dV$$

integrando ao longo do volume de controle,

$$\Sigma \vec{F}_{ext} = -\rho_1 V_1^2 A_1 + \rho V_2^2 A_2 + \frac{d}{dt} (\rho V \Psi)$$

$$\Sigma \vec{F}_{ext} = -\rho_1 V_1^2 A_1 + \rho V_2^2 A_2 + \rho V \frac{d\Psi}{dt} + \Psi \left(V \frac{d\rho}{dt} + \rho \frac{dV}{dt} \right)$$

mas, o balanço de forças que atuam no gás confinado no cilindro,

$$\Sigma \vec{F}_{ext} = -M_p a_p - M_g g + F_\mu$$

onde,

M_p = massa do pistão, dos anéis e de parte oscilante da biela

a_p = aceleração do pistão.

portanto,

$$\frac{dV}{dt} = \frac{1}{\rho \Psi} \left[-M_p a_p - M_g g + F_\mu + \rho_1 V_1^2 A_1 - \rho V_2^2 A_2 - \rho V \frac{d\Psi}{dt} - \Psi V \frac{d\rho}{dt} \right]$$

trocando a variável independente,

$$\frac{dV}{d\theta} = \frac{1}{\rho \Psi \omega} \left[-M_p a_p - M_g g + F_\mu + \dot{m}_1 V_1 - \dot{m}_2 V_2 - \rho V \omega \frac{d\Psi}{d\theta} - \Psi V \omega \frac{d\rho}{d\theta} \right]$$

Equação de Estado: (Gás Ideal)

$$P = \rho R T$$

tirando o logaritmo e derivando em relação ao ângulo do virabrequim,

$$\frac{1}{P} \frac{dP}{d\theta} = \frac{1}{\rho} \frac{d\rho}{d\theta} + \frac{1}{R} \frac{dR}{d\theta} + \frac{1}{T} \frac{dT}{d\theta}$$

$$\frac{dP}{d\theta} = \frac{P}{\rho} \frac{d\rho}{d\theta} + \frac{P}{R} \frac{dR}{d\theta} + \frac{P}{T} \frac{dT}{d\theta}$$

onde,

$$R = \sum_{i=1}^N C_i R_i$$

$$\frac{dR}{d\theta} = \sum_{i=1}^N R_i \frac{dC_i}{d\theta}$$

D.2 DEPURAÇÃO QUEIMA NO CILINDRO UMA ZONA

Conservação da massa:

$$\frac{d\rho}{d\theta} = \frac{1}{\omega \Psi} \left(\dot{m}_1 + \dot{m}_{inj} - \dot{m}_2 - \rho \omega \frac{d\Psi}{d\theta} \right)$$

Eliminando o fluxo de massa, temos,

$$\frac{d\rho}{d\theta} = \frac{1}{\Psi} \left(-\rho \frac{d\Psi}{d\theta} \right)$$

Conservação da energia:

$$\frac{dT}{d\theta} = \frac{(\gamma-1)}{\rho \Psi \omega R} \left[-\rho \Psi \omega \left(V \frac{dV}{d\theta} \right) - \omega \left(h + \frac{V^2}{2} - \frac{P}{\rho} \right) \left(\Psi \frac{d\rho}{d\theta} + \rho \frac{d\Psi}{d\theta} \right) + \omega \frac{\delta Q_f}{d\theta} + \omega \frac{\delta Q_i}{d\theta} - \omega P \frac{d\Psi}{d\theta} - \omega \frac{\delta W_\mu}{d\theta} + \dot{m}_1 \left(h_1 + \frac{V_1^2}{2} \right) + \dot{m}_{inj} h_f - \dot{m}_2 \left(h + \frac{V_2^2}{2} \right) \right]$$

Desprezando alguns termos e eliminando o fluxo de massa através do VC, temos,

$$\frac{dT}{d\theta} = \frac{(\gamma-1)}{\rho \Psi R} \left[- \left(c_p T - \frac{P}{\rho} \right) \left(\Psi \frac{d\rho}{d\theta} + \rho \frac{d\Psi}{d\theta} \right) + \frac{\delta Q_f}{d\theta} - P \frac{d\Psi}{d\theta} \right]$$

Equação de estado:

$$\frac{1}{P} \frac{dP}{d\theta} = \frac{1}{\rho} \frac{d\rho}{d\theta} + \frac{1}{T} \frac{dT}{d\theta}$$

$$\gamma = \frac{c_p}{c_v} = \frac{c_p}{c_p - R}$$

$$c_p = \frac{\gamma}{(\gamma-1)} R$$

Portanto,

$$\frac{dT}{d\theta} = \frac{(\gamma-1)}{\rho \Psi R} \left[- \left(\frac{\gamma}{(\gamma-1)} RT - \frac{P}{\rho} \right) \left(\Psi \frac{d\rho}{d\theta} + \rho \frac{d\Psi}{d\theta} \right) + \frac{\delta Q_f}{d\theta} - P \frac{d\Psi}{d\theta} \right]$$

Mas, da equação da conservação da massa,

$$\Psi \frac{d\rho}{d\theta} + \rho \frac{d\Psi}{d\theta} = 0$$

Assim,

$$\frac{dT}{d\theta} = \frac{(\gamma-1)}{\rho \Psi R} \left[\frac{\delta Q_f}{d\theta} - P \frac{d\Psi}{d\theta} \right]$$

Usando a equação do gás ideal na forma diferencial, temos,

$$\frac{1}{T} \frac{dT}{d\theta} = \frac{1}{P} \frac{dP}{d\theta} - \frac{1}{\rho} \frac{d\rho}{d\theta}$$

$$\frac{1}{T} \frac{dT}{d\theta} = \frac{1}{P} \frac{dP}{d\theta} - \frac{1}{\rho} \frac{d\rho}{d\theta}$$

Combinando com a conservação de massa,

$$\frac{dT}{d\theta} = \frac{T}{P} \frac{dP}{d\theta} - \frac{T}{\rho} \frac{d\rho}{d\theta} = \frac{T}{P} \frac{dP}{d\theta} + \frac{T}{V} \frac{dV}{d\theta}$$

Assim,

$$\frac{T}{P} \frac{dP}{d\theta} + \frac{T}{V} \frac{dV}{d\theta} = \frac{(\gamma-1)}{\rho V R} \left[\frac{\delta Q_f}{d\theta} - P \frac{dV}{d\theta} \right]$$

Mas, da equação do gás ideal,

$$\rho R = \frac{P}{T}$$

Assim,

$$\frac{T}{P} \frac{dP}{d\theta} + \frac{T}{V} \frac{dV}{d\theta} = \frac{T(\gamma-1)}{P V} \left[\frac{\delta Q_f}{d\theta} - P \frac{dV}{d\theta} \right]$$

$$\frac{\delta Q_f}{d\theta} - P \frac{dV}{d\theta} = \frac{1}{(\gamma-1)} \left[V \frac{dP}{d\theta} + P \frac{dV}{d\theta} \right]$$

$$\frac{\delta Q_f}{d\theta} = \frac{1}{(\gamma-1)} \left[V \frac{dP}{d\theta} + P \frac{dV}{d\theta} \right] + P \frac{dV}{d\theta}$$

$$\frac{dT}{d\theta} = T \left[\frac{1}{P} \frac{dP}{d\theta} + \frac{1}{V} \frac{dV}{d\theta} \right]$$

$$\frac{d\rho}{d\theta} = \frac{1}{V} \left(-\rho \frac{dV}{d\theta} \right) = -\frac{\rho}{V} \frac{dV}{d\theta}$$

D.3 DEPURAÇÃO DE RAZÃO DE SANGRIA NO TURBOCOMPRESSOR

Balço de energia no turbocompressor:

$$\dot{m}_c c_{pc} (T_{2c} - T_{1c}) = \dot{m}_t c_{pt} (T_{2t} - T_{1t}) + (\dot{m}_c - \dot{m}_t) c_{pt} T_{1t}$$

Razão de temperaturas: (x = c ou t) $\tau_x = \frac{T_{2x}}{T_{1x}}$

$$\dot{m}_c c_{pc} T_{1c} (\tau_c - 1) = \dot{m}_t c_{pt} T_{1t} (\tau_t - 1) + (\dot{m}_c - \dot{m}_t) c_{pt} T_{1t}$$

Razão de bypass pela válvula de sangria da turbina: $\beta = \frac{(\dot{m}_c - \dot{m}_t)}{\dot{m}_c} = 1 - \frac{\dot{m}_t}{\dot{m}_c}$

$$c_{pc} T_{1c} (\tau_c - 1) = (1 - \beta) c_{pt} T_{1t} (\tau_t - 1) + \beta c_{pt} T_{1t}$$

$$c_{pc} T_{1c} (\tau_c - 1) = c_{pt} T_{1t} (\tau_t - 1) - \beta c_{pt} T_{1t} (\tau_t - 1) + \beta c_{pt} T_{1t}$$

$$\beta = \frac{c_{pc} T_{1c} (\tau_c - 1) - c_{pt} T_{1t} (\tau_t - 1)}{-c_{pt} T_{1t} (\tau_t - 1) + c_{pt} T_{1t}}$$

$$\beta = \frac{c_{pc} T_{1c} (\tau_c - 1) - c_{pt} T_{1t} (\tau_t - 1)}{c_{pt} T_{1t} (2 - \tau_t)}$$